

# AN EXPERIMENTAL INVESTIGATION OF THE USE OF INJECTION LOCKING ON THE REMODULATION AND FILTERING OF WDM CHANNELS

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## ABSTRACT

The use of the optical injection locking in WDM receivers is experimentally investigated. As the locking bandwidth is controlled by the injected optical power, only the WDM carrier within the locking range can induce the locking of the slave laser and, therefore, be properly transmitted. In addition, properties of remodulation and reshaping for the selected channel information can be achieved by combining the injection locking process with a feed-forward laser current scheme. It was observed that optically injected semiconductor lasers can be used as optical filters, offering up to 40 dB extinction ratio. Also, the same locked laser under feed-forward current operation can perform the remodulation of 940 Mb/s optical channels.

## 1. INTRODUCTION

Recently, multiplexing techniques [1] have been combined with different physical transmission media [2] as a way to provide a solution for the constantly increasing demands on the expansion of the communication system capacity. After the development of Erbium doped fiber amplifiers (EDFA), the application of multiplexing techniques to optical systems became commercially possible, allowing a more realistic exploration of the optical fiber potential bandwidth. In optical communications, it is adopted the wavelength division multiplexing (WDM), where several channels allocated at different optical wavelengths are combined and transmitted by a single optical fiber. The system can support high bit rates and offer, at same time, security and reliability. At the receiver end, the channels are separated when samples of the optical signal are filtered out by optical filters centered at the different channel carrier wavelengths. Unfortunately, the wide bandwidth of the optical filters is one of the causes for the restriction on the maximum number of transmission channels in WDM systems, that can be amplified by EDFAs [1].

Under optical injection locking (OIL), a slave laser (SL), operates at the same frequency as that of a master laser (ML). During the locking condition, the SL gain dynamics are altered, inducing the SL noise characteristics to be the same as those of the ML. The two lasers remain locked while the equivalent free-running (no injection) frequency difference between them is within the OIL range. Normally, the locking range is narrow ( $< 10$  GHz) and controlled by the amount of injected optical power. If several WDM optical carriers were simultaneously coupled into a SL, only the channel whose frequency is inside the locking range would induce the SL locked operation. In this way, the SL would reproduce the ML signal characteristics with high output power, while attenuating the other channels by optical

absorption. As a result, the OIL could be used in the optical filtering of WDM channels with the advantage of allowing closer channel spacing due to the narrow locking range.

Also, once inside the locking range, the SL output optical power varies almost linearly, depending on the equivalent free-running frequency difference between master and slave lasers (detuning). If semiconductor lasers are used, the optical power reaches a maximum for negative frequency detuning (ML frequency  $>$  SL frequency) and a minimum for positive detuning. Therefore, if the free-running frequency difference can be adequately controlled, it is possible to vary the level of the SL output optical power accordingly. As in semiconductor lasers the optical frequency is tuned by the electronic bias current, it would be possible alter the equivalent value of the SL free-running frequency to properly amplitude modulate the locked laser. In this way, the OIL process could also be employed in the remodulation and/or reshaping of WDM optical channels, as the incoming signal is substituted by the processed SL signal.

In this paper, the utilization of the optical injection locking technique in WDM receivers for filtering and remodulation/reshaping of optical channels is experimentally investigated. First, theoretical concepts concerning the optical injection locking process are studied. Following, the block diagram for the OIL filtering experimental set-up is presented. A feed-forward approach was implemented in the initial set-up to allow the remodulation/reshaping investigation to take place. Finally, the experimental results are presented and analyzed.

## 2. THEORY

Fig. 1 shows the a simplified block diagram of an OIL experiment [3]. The ML light is injected into the SL active region, after an isolator. The isolator prevents SL coupling into the ML. If the ML and SL optical frequencies are sufficiently close, that is, inside the locking range, the gain mechanisms of the SL are altered in such a way to force its optical frequency to be the same as that of the ML. The lasers are kept locked while the free-running frequency difference between ML and SL is inside a so called locking range. As mentioned before, under OIL conditions, both lasers operate at the same frequency, with the SL following all phase fluctuations of the ML.

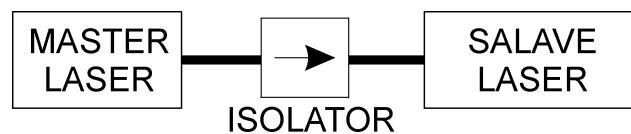


Figure 1. Block diagram for the OIL system.

The interaction of the electric field with the leasing medium can be evaluated by a set of coupled differential rate equations, describing the photon number, the phase and the carrier number. Assuming that the lasers are single mode and the system is locked, the SL rate equations can be expressed by [4-5]:

$$\frac{dI_s(t)}{dt} = \left[ \text{Re}[G(t)] - \frac{1}{\tau_p} \right] I_s(t) + \frac{2\sqrt{\eta I_s(t) I_m(t)}}{\tau_i} \cos \theta(t) + R \quad (1)$$

$$\frac{d\phi_s(t)}{dt} = \omega_s - \omega_m + \frac{1}{2} \text{Im}[G(t)] + \frac{1}{\tau_i} \sqrt{\frac{\eta I_m(t)}{I_s(t)}} \sin \theta(t) \quad (2)$$

$$\frac{dN(t)}{dt} = \frac{i(t)}{q} - \frac{N(t)}{\tau_s} - \text{Re}[G(t)] I_s(t) \quad (3)$$

where  $I_s(t)$  and  $I_m(t)$ ,  $\phi_s(t)$  and  $\phi_m(t)$ , and  $\omega_s$  and  $\omega_m$  are, respectively, the SL and ML photons numbers, phases (rad), and free-running angular frequencies (rad/s),  $\theta(t) = \phi_m(t) - \phi_s(t)$ ,  $G(t)$  is the complex gain ( $s^{-1}$ ),  $\tau_p$  is the photon life time (s),  $\eta$  is a term representing the coupling efficiency of the master laser injected signal,  $\tau_i$  is the round trip time of the electric field inside the laser structure (s),  $R$  is the spontaneous emission rate ( $s^{-1}$ ),  $N(t)$  is the carrier number,  $\tau_s$  is the carrier lifetime (s),  $i(t)$  is the laser current (A),  $q$  is the electron charge (C) and  $d$  is the active layer thickness (m). In order to simplify the coupled non-homogenous differential equations (1) to (3), the time-dependent terms can be linearized and solved using the approximation of small perturbations around the steady state values. In this context, a given variable  $A(t)$  can be expressed as  $A(t) = A_o + \hat{a}(t)$ , where  $A_o$  is the  $A(t)$  stationary value and  $\hat{a}(t)$  (where  $\hat{a}(t) \ll A_o$ ) is the perturbation around the stationary value  $A_o$ . The complex SL gain  $G(t)$ , within a first order approximation, can be written as [4]:

$$G(t) = G_o + G_I \hat{I}_s(t) + G_N (1 + j\alpha) [\Delta N_o + \hat{n}(t)] \quad (4)$$

where  $\Delta N_o = N_o - \tilde{N}_o$ ,  $N_o$  is the SL carrier number under injection,  $\tilde{N}_o$  represents the free-running steady state carrier number,  $G_o$  is the free-running gain per unit of time ( $s^{-1}$ ),  $G_N = \delta G / \delta N$  is differential gain related to the carrier number ( $s^{-1}$ ),  $G_I = \delta G / \delta I$  is differential gain ( $s^{-1}$ ) related to the photon number, and  $\alpha$  is known as the linewidth enhancement factor and accounts for the phase-amplitude coupling of the electric field. The  $\hat{I}_s(t)$  and  $\hat{n}(t)$  terms represent the oscillations (perturbations) around the stationary values for photon number and carrier number, respectively.

For the purposes of this paper, the stationary solutions of the linearized differential equations are sufficient. After a proper mathematical manipulation, the stationary solutions for the OIL process are given by [5-6]:

$$\Delta G = G_N \Delta N_o = -\frac{v_g}{L} \sqrt{\frac{\eta I_{mo}}{I_{so}}} \cos \theta_o \quad (5)$$

$$\Delta \omega = \omega_m - \omega_s = -\frac{v_g}{2L} \sqrt{\frac{\eta I_{mo}}{I_{so}}} (\sin \theta_o - \alpha \cos \theta_o) \quad (6)$$

$$I_{so} = \frac{G_o \tilde{I}_{so} - \frac{\Delta N_o}{\tau_s}}{G_o + \Delta G} \quad (7)$$

where  $I_{so}$  is the SL steady state photon number under injection,  $I_{mo}$  is the ML steady state photon number,  $\tilde{I}_{so}$  is the free-running SL steady state photon number,  $v_g$  is the group velocity (m/s),  $L$  is the SL effective cavity length (m), and  $\theta_o$  is the steady state phase difference between ML and SL (rad).

Equation (6) is of particular interest for the injection locking process. It defines the injection locking bandwidth of the OIL system [6]. It is possible to observe that the width of the locking range is related to the amount of injected optical power. Therefore, for a given injection level, locking can only be achieved if the equivalent free-running frequency difference between the ML and SL falls within this range. Therefore, in WDM systems, any injected optical channel whose frequency lays outside the locking bandwidth would cause no changes in the dynamics of the SL, being mainly attenuated by absorption. On the other hand, only the channel wavelength that causes laser locking would be properly transmitted and photodetected.

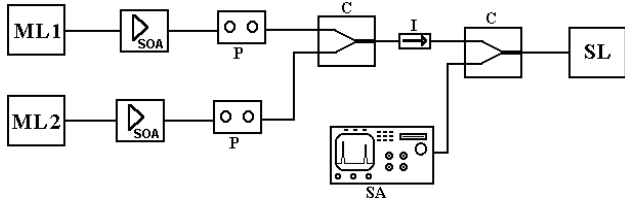
As it can be seen in (5) and (7), the SL gain ( $\Delta G$ ) and photon number ( $I_{so}$ ) are also susceptible to changes due to oscillations in the injected ML photon number. Hence, it is quite possible that the SL output signal presents the same oscillating nature as that of the ML signal. By recalling that the electronic SL bias current can also alter the photon number (through  $G_o$ ), a proper combination of effects would allow to maximize or minimize the oscillation contents of the SL signal. If the oscillations in the ML signal are seen as intensity modulation, the SL laser could perform the remodulation, reshaping, and/or erasing of the information in a given WDM optical channel.

### 3. EXPERIMENTAL RESULTS

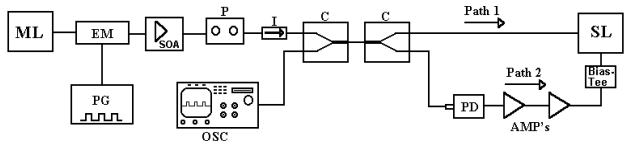
Fig. 2 shows the block diagram for the OIL filtering experimental set-up. The signals from two semiconductor lasers, ML1 and ML2, are amplified and simultaneously coupled into active region of the SL, after two 3dB fiber couplers. In a single facet arrangement, the light emitted by SL is coupled into an optical spectrum analyzer. Isolators prevent the SL light and retro-reflections to be coupled into ML1 and ML2. Polarization controllers are used to accentuate the polarization matching of different wavefronts. By assuming that the SL is selected to filter out the ML2 signal, the coupling of the two ML signals results in the locking of the SL in relation to ML1. Thus, the ML2 photons have practically no influence in the SL gain mechanism changes and the stimulated emissions in this frequency are rather small. As a result, the photons at the ML2 frequency tend to spread or be absorbed inside the SL active region. Consequently, only a small number of ML2 photons are able to leave the SL cavity and be detected.

Fig. 3 shows the experimental block diagram for the investigation of the OIL remodulation, reshaping and erasing properties. In the experiment, the ML signal is externally modulated, amplified by a semiconductor optical amplifier (SOA), and divided into two different optical paths. One part of the ML signal is injected into the SL (path 1). The other is

photodetected (path 2). The AC photocurrent is amplified and combined with the SL bias current. The process described for path 2 is responsible for the current control of the SL. In other words, a sample of the modulated ML optical signal is fed-forward into the SL via the bias current. If the two path lengths are closely matched, the modulated ML signal is electronically and optically coupled into the SL and, as previously described, the oscillations caused by the intensity modulation produce simultaneous gain and photon number variations. Hence, with an adequate phase combination for the optical and electronic signals, it is possible to obtain remodulation and/or reshaping of the optical carrier information. An optical oscilloscope was used to analyze the signals before and after the SL.

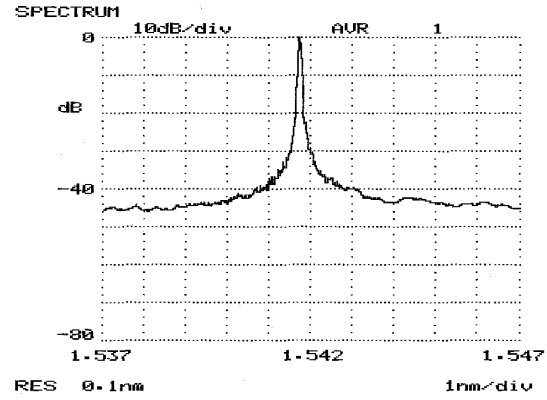


**Figure 2.** Block diagram for optical filtering experiment using the OIL process. SOA: semiconductor optical amplifier; P: polarization controller; I: isolator; C: fiber coupler; SA: optical spectrum analyzer.

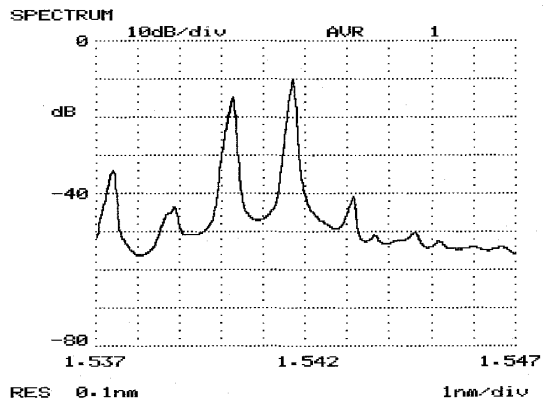


**Figure 3.** Block diagram for the OIL remodulation experiment. PG: pulse generator; EM: external modulator; P: polarization controller; I: isolator; C: fiber coupler; PD: photodetector; OSC: optical signal analyzer (oscilloscope); SOA: semiconductor optical amplifier.

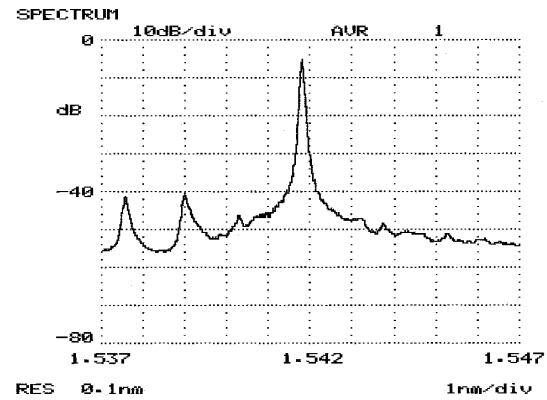
To investigate the injection locking effect, the set-up shown in Fig. 2 was adapted for the measurement of the optical spectra of the ML, the SL, and the locked SL. The results are presented in Fig. 4. The master laser (ML1 only) was a tunable 1550 nm external cavity semiconductor laser (Photonics) operating at 60 mA bias current. The SL was a base-mounted Fabry-Perot (FP) semiconductor laser, operating at 60 mA. By referring back to Fig. 2, at the SL end, it was necessary to use bulk optics to allow access to the only available SL facet. Hence, the ML1 signal was collimated after the second optical fiber coupler and coupled into the SL cavity. Reflections and the SL emission followed the same optical path, but in the opposite direction. To measure Fig. 4(a), the optical spectrum analyzer was connected after the isolator in Fig. 2. The measurements of Fig. 4(b) and (c) were obtained as in the original set-up. The ML1, Fig. 4(a), was mechanically tuned to closely match one of the SL mode wavelengths, as shown in Fig. 4(b). By controlling the SL temperature, locking happened at 1541.75 nm, Fig. 4(c). It can be seen that the SL multimode characteristics shown in Fig. 4(b) disappear once the SL operates in the locked mode. As expected from the injection locking process, the SL phase noise content is suppressed and its spectrum becomes similar to that of the master laser, as in Fig. 4(a) and (c).



(a)



(b)



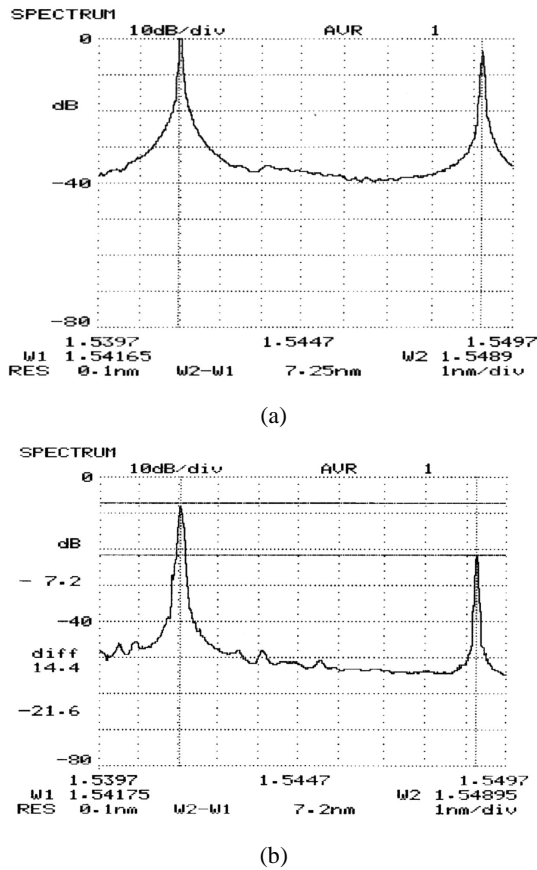
(c)

**Figure 4.** Optical spectrum of master laser (a), slave laser (b), and locked slave laser (c).

In the OIL filtering experiment, ML1 and SL are the same as above. The ML2 was a non-tunable external cavity semiconductor laser (E-Tek Ind.) operating at 1548.95 nm and 120 mA bias current. The first coupler in Fig. 2 is responsible by the combination of the ML1 and ML2 signals. Fig 5(a) shows the ML1 and ML2 optical spectra, measured after isolator output. It is possible to observe that the difference between the ML1 and

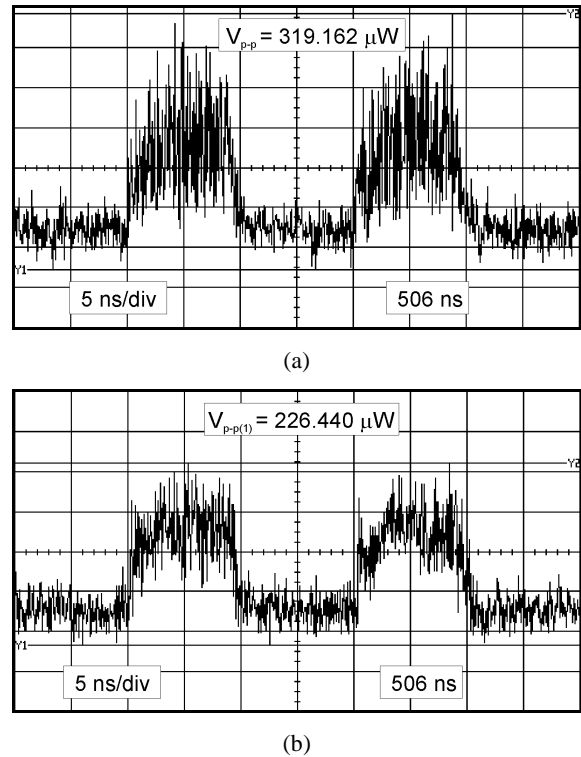
ML2 power levels is around 4 dB. Fig. 5(b) shows the measured spectra when the spectrum analyzer is used as in Fig. 2. The OIL filtering process increases the power level difference to around 14 dB. Nevertheless, due to the discrete components used in the experiment, a strong optical reflection in the fiber-air interface was observed after the second coupler. Additional measurements verified that the ML2 power level could drop further 25 dB if the reflection effect were eliminated.

To perform the experiment represented by Fig. 3, the set-up was modified and LM2 was removed. All operational characteristics of the ML1 and SL were kept the same. An external modulator (Sumitomo Inc.) modulated the ML1 signal, using a HP 8131A signal generator. A 300kHz-6GHz photodetector (HP 83411A) was used to detect a sample of ML1 optical signal. The AC current coupled into the SL was 3 mA<sub>RMS</sub>. The SL bias current was 60 mA. Fig. 6 shows the plots obtained from an optical oscilloscope (HP 83480), where the effect of the OIL over the direct modulation of the SL current by the electronic ML sampled signal at 100 Mb/s is observed. It is possible to verify that, under locking, Fig. 6(b), the modulated SL signal becomes less noisy than the signal with no injection, Fig. 6(a). The reason for this behavior comes from the fact that the SL, under OIL, acquires the ML noise characteristics. As the ML is an external cavity laser, its low phase noise contents are passed to the noisy FP laser when locking is achieved.



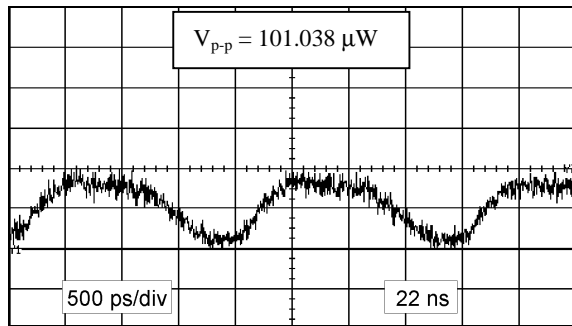
**Figure 5.** Optical filtering using OIL process, (a) no injection (b) with injection.

Fig. 7 shows the results for the information processing experiment, when the bit rate is set at 940 Mb/s. Fig 7(a) shows the photodetected SL signal when no locking is present and the SL is modulated by the feed-forward current. Once that SL is locked, the phase adjustment between the optical and electronic injection signals generate two distinct effects. In the first one, both electronic and optical signals are in phase, Fig. 7(b), causing the peak-to-peak value of the SL modulated signal to increase from 101  $\mu$ W, Fig 7(a), to 162  $\mu$ W. Although the remodulated pulse shape is different from that of the pulse without injection, the pulse central region in Fig. 7(b) allows a clear distinction between logic levels so that clock recovery would still be possible. Fig 7(c) presents the erasing effect, when the signals are near to a 180-degree detuning. It is possible to observe that the differences between the maximum and minimum levels decrease considerably. By observing the peak-to-peak value in Fig. 7(c), it is almost impossible to distinguish between logic levels. This response would presumably induce logical receivers to interpret the data stream as composed only by low logic levels.

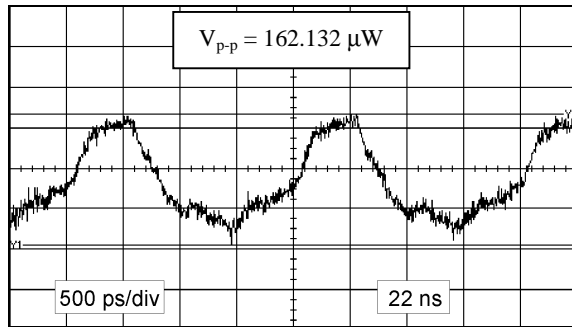


**Figure 6.** SL under direct modulation of the electronic sample of the ML signal without (a) and with locking (b).

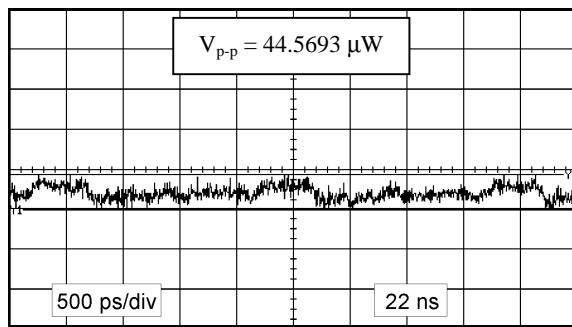
It is believed that these results can be optimized if a better quality SL is used in the experiment. The available SL had a low quantum efficiency and a poor direct modulation response. Also, its complex structure made optical coupling difficult, so that ML signal amplification was necessary. The adjustment of phase difference between optical and electronic signals was also difficult. The phase control was obtained from fine tuning of the operation frequency (470 MHz or 940 Mb/s) in the pulse generator. As a result, the signal remodulation was poor. Despite the problems, it was possible to demonstrate the feasibility of the process.



(a)



(b)



(c)

**Figure 7.** SL output signal with feed-forward: without locking (a), under remodulation (b), and with erasing (c). Oscilloscope mode: average = 16.

## 4. CONCLUSIONS

In this paper, the utilization of the optical injection locking technique for filtering and remodulation/reshaping of optical channels was experimentally investigated. It was observed that the OIL filtering has the potential to suppress undesirable optical carriers by up to 40 dB. Also, it was demonstrated that the OIL can be used in the remodulation (amplification and erasing) of the information in optical carriers. In the amplification process, the signal level was almost doubled. In the erasing application, it dropped to less than 40 % of the original value. In spite of the deficiencies observed in the experimental set-up, the results were within the expected. A theory review is necessary to compensate for the linearization adopted in Section II, which is only valid within certain limits. It is expected that a more detailed analysis

of OIL process could help in the development of more sophisticated electronic projects to improve the SL signal coupling.

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